

# On the Stability of Kähler-Einstein Metrics

Xianzhe Dai\*      Xiaodong Wang†      Guofang Wei‡

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## Abstract

Using  $\text{spin}^c$  structure we prove that Kähler-Einstein metrics with nonpositive scalar curvature are stable (in the direction of changes in conformal structures) as the critical points of the total scalar curvature functional. Moreover if all infinitesimal complex deformation of the complex structure are integrable, then the Kähler-Einstein metric is a local maximal of the Yamabe invariant, and its volume is a local minimum among all metrics with scalar curvature bigger or equal to the scalar curvature of the Kähler-Einstein metric.

## 1 Introduction

Stability issue comes up naturally in variational problems. One of the most important geometric variational problems is that of the total scalar curvature functional. Following [Bes87, Page 132] we call an Einstein metric stable if the second variation of the total scalar curvature functional is nonnegative in the direction of changes in conformal structures. By the well-known formula, this is to say,

$$\langle \nabla^* \nabla h - 2\overset{\circ}{R}h, h \rangle \geq 0 \tag{1.1}$$

for any trace-free and divergence-free symmetric two tensor  $h$ . Here  $\overset{\circ}{R}h$  denotes the natural action of the curvature tensor on the symmetric tensors [Bes87]. The operator appearing in (1.1) is closely related to the Lichnerowicz Laplacian. In fact, one has

$$\mathcal{L}_g h = \nabla^* \nabla h - 2\overset{\circ}{R}h + 2\text{Ric} \circ h. \tag{1.2}$$

The two thus coincide for Ricci flat metrics.

In [DWW04], we studied the stability of compact Ricci flat manifolds. An essential ingredient there is the use of spin structure and parallel spinors. In fact, our result should really be viewed as the stability result for compact Riemannian manifolds with nonzero parallel spinor. By [Wa89], [H74-1], this class of manifolds essentially coincides with that of special holonomy, namely, the Calabi-Yau manifolds, hyperKähler manifolds,  $\text{spin}(7)$  manifolds and  $G_2$  manifolds.

In this paper, we use  $\text{spin}^c$  structure to generalize our previous result to manifolds with nonzero parallel  $\text{spin}^c$  spinor. Since the existence of nonzero parallel spinor implies that the metric is necessarily Ricci flat, our motivation here is to extend our previous method to deal with nonzero scalar curvature and we found  $\text{spin}^c$  to be a good framework to work with.

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\*Math Dept, UCSB, Santa Barbara, CA 93106 Email: dai@math.ucsb.edu. Partially supported by NSF Grant # DMS-0405890

†Math Dept, Michigan State University, East Lansing, MI 48824 Email: xwang@math.msu.edu.

‡Math Dept. UCSB. Email: wei@math.ucsb.edu. Partially supported by NSF Grant # DMS-0204187.

**Theorem 1.1** *If a compact Einstein manifold  $(M, g)$  with nonpositive scalar curvature admits a nonzero parallel  $\text{spin}^c$  spinor, then it is stable.*

It turns out that manifolds admitting a nonzero parallel  $\text{spin}^c$  spinor are more or less classified [Mo97]. Namely a simply connected manifold has a nonzero parallel  $\text{spin}^c$  spinor if and only if the manifold is the product of a Kähler manifold and a manifold with parallel spinor. Moreover, the  $\text{spin}^c$  structure is the product of the canonical  $\text{spin}^c$  structure on the Kähler manifold with the spin structure on the other factor. The case with parallel spinor being treated in [DWW04], we only need to deal with the Kähler case.

For manifolds with nonzero parallel  $\text{spin}^c$  spinor, we derive a Bochner type formula relating the operator  $\nabla^*\nabla - 2\overset{\circ}{R}$  to the square of a twisted Dirac operator. The difference, which is expressed in terms of the curvatures, can be shown to be nonnegative under our assumption. In fact, we prove that the operator  $\nabla^*\nabla - 2\overset{\circ}{R}$  is positive semi-definite for Kähler manifolds with nonpositive Ricci curvature.

The operator  $\nabla^*\nabla - 2\overset{\circ}{R}$  (or the Lichnerowicz Laplacian) seems to have a knack for appearing in geometric variational problems. Besides the total scalar curvature functional, or equivalently, the Yamabe functional if one normalizes the volume,

$$Y(g) = \frac{\int_M S_g dV_g}{\text{Vol}(g)^{1-\frac{2}{n}}},$$

there is also the first eigenvalue  $\lambda(g)$  of conformal Laplacian considered in [DWW04], and the  $L^{n/2}$  norm of scalar curvature [BCG91]

$$K(g) = \int_M |S_g|^{n/2} dV_g.$$

Using these functionals, we can then deduce a number of interesting consequences.

**Theorem 1.2** *Let  $(N, g_0, J_0)$  be a compact Kähler-Einstein manifold with nonpositive scalar curvature. Suppose all infinitesimal complex deformations of  $J_0$  are integrable. Then there exists a neighborhood  $\mathcal{U}$  of  $g_0$  in the space of smooth Riemannian metrics on  $N$  such that for any metric  $g \in \mathcal{U}$  with scalar curvature  $S_g \geq S_{g_0}$*

$$\text{Vol}(N, g) \geq \text{Vol}(N, g_0)$$

*and equality holds iff  $g$  is a Kähler-Einstein metric with negative scalar curvature.*

In the case of zero scalar curvature, the integrability condition is automatic by the Bogomolov-Tian-Todorov theorem [Bo78], [T86], [To89].

**Theorem 1.3** *Let  $(N, g_0, J_0)$  be a compact Kähler-Einstein manifold with nonpositive scalar curvature. Suppose all infinitesimal complex deformations of  $J_0$  are integrable. Then  $g_0$  is a local maximum of the Yamabe invariant.*

There are many examples satisfying the assumptions in the theorems above. For example, the hypersurfaces of large enough degree in a complex projective space. In fact we do not know any examples of Kähler-Einstein manifolds of nonpositive scalar curvature which do not satisfy the integrability condition for the complex structure and we believe that they all satisfy the integrability condition, just as Calabi-Yau manifolds by virtue of the Bogomolov-Tian-Todorov theorem [Bo78], [T86], [To89].

The study of the Yamabe constant, also called Schoen's  $\sigma$  invariant, has attracted a lot of attention lately, Cf. [Le99], [BN04]. This is motivated by a conjecture of Schoen [Sch89], which says that the standard metric for manifolds with constant sectional curvature realizes the Yamabe constant. In other words, the standard metric is a global maximum for the Yamabe invariant (they are called the supreme Einstein metrics in [Le99]). In view of the results of [BCG95] for real hyperbolic spaces and of [Le95] for Kähler-Einstein surfaces, it is tempting to conjecture the same for more general class of manifolds such as compact locally symmetric spaces or even Kähler-Einstein manifolds. Unfortunately, this is not true in higher dimensions, as a compact simply connected manifold of dimension greater than or equal to 5 must have nonnegative Yamabe constant [Pe00].

There has been a lot of work recently concerning the stability of Ricci flow [GIK02], [Se04], [Ch05], see also [CHI04]. The general question can be phrased as follows. If  $g_0$  is a metric such that the (renormalized) Ricci flow  $g(t)$  starting from  $g_0$  converges, is it true that the (renormalized) Ricci flow  $\tilde{g}(t)$  starting from all metrics  $\tilde{g}_0$  that are sufficiently close to  $g_0$  also converges? Using the result of Natasa Sesum [Se04], we derive

**Theorem 1.4** *Let  $(N, g_0, J_0)$  be a compact Kähler-Einstein manifold with nonpositive scalar curvature. Suppose all infinitesimal complex deformations of  $J_0$  are integrable. Then the Ricci flow starting from any Riemannian metric sufficient close to  $g_0$  converges exponentially to a Kähler-Einstein metric diffeomorphic to  $g_0$ .*

The difference between this theorem and the well known result for Kähler-Ricci flow on Kähler-Einstein manifolds with nonpositive first Chern class [Cao85] is that the Ricci flow here starts with any metric, rather than in a given Kähler class. On the other hand, the result of [Cao85] is a global result.

This paper is organized as follows. We discuss  $\text{spin}^c$  parallel spinor and related Bochner type formula in the next section, and prove the infinitesimal stability result. In Section 3, we discuss the local stability results and applications. In the final section, relevant results from Kodaira-Spencer theory are recalled. We also elaborate more on the examples and make some remarks.

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## 2 $\text{Spin}^c$ parallel spinor

We now assume  $(M, g)$  is a compact Riemannian manifold with a  $\text{spin}^c$  structure. Thus,  $w_2(M) \equiv c$ , where  $c \in H^2(M, \mathbb{Z})$  is the canonical class of the  $\text{spin}^c$  structure. Let  $\mathcal{S}^c \rightarrow M$  denote the  $\text{spin}^c$  spinor bundle and  $L \rightarrow M$  the complex line bundle with  $c_1(L) = c$ . Then  $\mathcal{S}^c = \mathcal{S} \otimes L^{1/2}$ , where the spinor bundle  $\mathcal{S}$  may not exist globally; similarly for the square root of the line bundle. An excellent reference on spin geometry is Lawson and Michelsohn [LM89].

Let  $E \rightarrow M$  be a vector bundle with a connection. The curvature is defined as

$$R_{XY} = -\nabla_X \nabla_Y + \nabla_Y \nabla_X + \nabla_{[X, Y]}. \quad (2.1)$$

If  $M$  is a Riemannian manifold, then for the Levi-Civita connection on  $TM$ , we have  $R(X, Y, Z, W) = \langle R_{XY}Z, W \rangle$ . We often work with an orthonormal frame  $\{e_1, \dots, e_n\}$  and its dual frame  $\{e^1, \dots, e^n\}$ . Set  $R_{ijkl} = R(e_i, e_j, e_k, e_l)$ .

The spinor bundle  $\mathcal{S}$ , which may exist only locally, has a natural connection induced by the Levi-Civita connection on  $TM$ . For a spinor  $\sigma$ , we have

$$R_{XY}\sigma = \frac{1}{4}R(X, Y, e_i, e_j)e_i e_j \cdot \sigma. \quad (2.2)$$

Given a unitary connection  $\nabla^L$  on  $L$ , we then obtain a Clifford connection  $\nabla^c$  on  $\mathcal{S}^c$ . In fact,  $\nabla^c = \nabla \otimes 1 + 1 \otimes \nabla^{L^{1/2}}$  is the tensor product connection for  $\mathcal{S}^c$ . Therefore, for a  $\text{spin}^c$  spinor  $\sigma$ ,

$$R_{XY}\sigma = \frac{1}{4}R(X, Y, e_i, e_j)e_i e_j \cdot \sigma - \frac{1}{2}F(X, Y)\sigma. \quad (2.3)$$

Here  $F$  is the curvature form of  $\nabla^L$ .

If  $\sigma_0$  is a parallel  $\text{spin}^c$  spinor, i.e.,  $\sigma_0$  is a section of  $\mathcal{S}^c$  such that  $\nabla_X^c \sigma_0 = 0$  for all  $X$ , then  $R_{XY}\sigma_0 = 0$ . Hence we have

$$R_{klij}e_i e_j \cdot \sigma_0 = 2F_{kl}\sigma_0. \quad (2.4)$$

**Lemma 2.1** *If  $\sigma_0$  is a parallel  $\text{spin}^c$  spinor, then*

$$R_{kl}e_l \cdot \sigma_0 = F_{kl}e_l \cdot \sigma_0.$$

**Proof:** From (2.4) we have

$$R_{klij}e_l e_i e_j \cdot \sigma_0 = 2F_{kl}e_l \cdot \sigma_0.$$

But

$$\begin{aligned} R_{klij}e_l e_i e_j &= \frac{1}{3} \sum_{l, i, j \text{ distinct}} (R_{klij} + R_{kijl} + R_{kjl i})e_l e_i e_j \\ &\quad + \sum_{i, j} R_{kji} e_j e_i e_j + \sum_{i, j} R_{kii} e_i e_i e_j \\ &= 2R_{ki}e_i. \end{aligned}$$

Here we have used the symmetries of Riemann curvature tensor, including the first Bianchi identity. Hence,

$$R_{kl}e_l \cdot \sigma_0 = F_{kl}e_l \cdot \sigma_0$$

as claimed. ■

From now on, we assume  $M$  has a parallel  $\text{spin}^c$  spinor  $\sigma_0 \neq 0$ , which, without loss of generality, is normalized to be of unit length. We define a linear map  $\Phi : S^2(M) \rightarrow \mathcal{S}^c \otimes T^*M$  by

$$\Phi(h) = h_{ij}e_i \cdot \sigma_0 \otimes e^j. \quad (2.5)$$

It is easy to check that the definition is independent of the choice of the orthonormal frame  $\{e_1, \dots, e_n\}$ .

**Lemma 2.2** *The map  $\Phi$  satisfies the following properties:*

1.  $\text{Re} \langle \Phi(h), \Phi(\tilde{h}) \rangle = \langle h, \tilde{h} \rangle$ ,
2.  $\nabla_X \Phi(h) = \Phi(\nabla_X h)$ .

Here  $\text{Re}$  denotes the real part.

**Proof:** We compute

$$\begin{aligned} \langle \Phi(h), \Phi(\tilde{h}) \rangle &= h_{ij}\tilde{h}_{kl} \langle e_i \cdot \sigma_0 \otimes e^j, e_k \cdot \sigma_0 \otimes e^l \rangle \\ &= h_{il}\tilde{h}_{kl} \langle e_i \cdot \sigma_0, e_k \cdot \sigma_0 \rangle \\ &= -h_{il}\tilde{h}_{kl} \langle \sigma_0, e_i e_k \cdot \sigma_0 \rangle \\ &= h_{kl}\tilde{h}_{kl} - \sum_{i \neq k} h_{il}\tilde{h}_{kl} \langle \sigma_0, e_i e_k \cdot \sigma_0 \rangle. \end{aligned}$$

Now, for  $i \neq k$ ,

$$\langle \sigma_0, e_i e_k \cdot \sigma_0 \rangle = \langle e_k e_i \cdot \sigma_0, \sigma_0 \rangle = -\langle e_i e_k \cdot \sigma_0, \sigma_0 \rangle = -\overline{\langle \sigma_0, e_i e_k \cdot \sigma_0 \rangle}.$$

That is,  $\langle \sigma_0, e_i e_k \cdot \sigma_0 \rangle$  is purely imaginary. Taking the real part of the previous equation proves the first assertion. To prove the second one, we choose our orthonormal frame such that  $\nabla e_i = 0$  at  $p$  and compute at  $p$

$$\begin{aligned} \nabla_X \Phi(h) &= X h_{ij} e_i \cdot \sigma_0 \otimes e^j \\ &= \nabla_X h(e_i, e_j) e_i \cdot \sigma_0 \otimes e^j \\ &= \Phi(\nabla_X h). \end{aligned}$$

■

The following interesting Bochner type formula plays an important role here.

**Lemma 2.3** *Let  $h$  be a symmetric 2-tensor on  $M$ . Then*

$$\mathcal{D}^* \mathcal{D} \Phi(h) = \Phi(\nabla^* \nabla h - 2\overset{\circ}{R}h + R^L(h) + \text{Ric}(h)). \quad (2.6)$$

Here  $R^L(h)_{ij} = h_{ip} F_{jp}$  and  $\text{Ric}(h)_{ij} = R_{ip} h_{jp}$ .

**Remark** Note that here we have implicitly extended our map  $\Phi$  to general (nonsymmetric) 2-tensors with complex coefficients.

**Proof:** Choose an orthonormal frame  $\{e_1, \dots, e_n\}$  near a point  $p$  such that  $\nabla e_i = 0$  at  $p$ . We compute at  $p$

$$\begin{aligned} \mathcal{D}^* \mathcal{D} \Phi(h) &= \nabla_{e_k} \nabla_{e_l} h(e_i, e_j) e_k e_l e_i \cdot \sigma_0 \otimes e^j \\ &= -\nabla_{e_k} \nabla_{e_k} h(e_i, e_j) e_i \cdot \sigma_0 \otimes e^j - \frac{1}{2} R_{e_k e_l} h(e_i, e_j) e_k e_l e_i \cdot \sigma_0 \otimes e^j \\ &= \Phi(\nabla^* \nabla h) + \frac{1}{2} R_{kljp} h_{ip} e_k e_l e_i \cdot \sigma_0 \otimes e^j + \frac{1}{2} R_{klip} h_{pj} e_k e_l e_i \cdot \sigma_0 \otimes e^j. \end{aligned}$$

By using twice the Clifford relation  $e_i e_j + e_j e_i = -2\delta_{ij}$  we have

$$\begin{aligned} \frac{1}{2} R_{kljp} h_{ip} e_k e_l e_i \cdot \sigma_0 &= \frac{1}{2} R_{kljp} h_{ip} e_i e_k e_l \cdot \sigma_0 + R_{kljp} h_{kp} e_l \cdot \sigma_0 - R_{kljp} h_{lp} e_k \cdot \sigma_0 \\ &= F_{jp} h_{ip} e_i \cdot \sigma_0 - 2(\overset{\circ}{R}h)_{kj} e_k \cdot \sigma_0. \end{aligned}$$

Here the last equality uses (2.4). On the other hand,

$$\frac{1}{2} R_{klip} h_{pj} e_k e_l e_i \cdot \sigma_0 = \frac{1}{2} h_{pj} (R_{klip} e_k e_l e_i \cdot \sigma_0) = R_{lp} h_{pj} e_l \cdot \sigma_0.$$

Putting these equations together we obtain our lemma. ■

The existence of a parallel  $\text{spin}^c$  spinor on a compact simply connected manifold implies that the manifold is the product of a Kähler manifold with a manifold with parallel spinor [Mo97]. Moreover, the  $\text{spin}^c$  structure is the product of the canonical  $\text{spin}^c$  structure on the Kähler manifold with the spin structure on the other factor.

We now assume that  $(M, g)$  is a compact Kähler manifold of real dimension  $n = 2m$ . Let  $J$  be the parallel almost complex structure and  $\omega = g(J \cdot, \cdot)$  the Kähler form. The complexified tangent bundle decomposes as

$$TM \otimes \mathbb{C} = T^{1,0}M \oplus T^{0,1}M.$$

The canonical  $\text{spin}^c$  structure is given by the anti-canonical line bundle  $L = K^{-1} = \Lambda^m(T^{1,0}(M))$ . It has a canonical holomorphic connection induced from the Levi-Civita Connection and the curvature form  $F = -\sqrt{-1}\rho$ , where  $\rho = \text{Ric}(J\cdot, \cdot)$  is the Ricci form.

The spinor bundle  $\mathcal{S}^c(M) = \mathcal{S}_+^c(M) \oplus \mathcal{S}_-^c(M)$  with

$$\begin{aligned}\mathcal{S}_+^c(M) &= \bigoplus_{k \text{ even}} \Lambda^{0,k}(M), \\ \mathcal{S}_-^c(M) &= \bigoplus_{k \text{ odd}} \Lambda^{0,k}(M).\end{aligned}$$

The Clifford multiplication is defined by

$$v \cdot = \sqrt{2}(v^{0,1} \wedge -v^{0,1} \lrcorner).$$

Here  $v^{0,1} \lrcorner$  denotes the contraction using the Hermitian metric. The parallel spinor  $\sigma \in C^\infty(\mathcal{S}_+^c(M))$  can be taken as the function which is identically 1.

**Remark** We would like to remark that the  $\text{spin}^c$  structure and related Bochner type formula are very useful in other context, such as symplectic manifolds. Given a symplectic manifold  $(M, \omega)$  of dimension  $2m$ , we take an almost complex structure  $J$  compatible with the symplectic form  $\omega$ . This gives rise to a Riemannian metric  $g = \omega(\cdot, J\cdot)$ . Then the formulation above in the Kähler case works perfectly well in this generalized setting and defines a natural  $\text{spin}^c$  structure on  $M$ . The Levi-Civita connection  $\nabla$  of  $g$  induces a natural Hermitian connection  $A$  on  $\Lambda^m(T^{1,0}(M))$  and hence a connection on the spinor bundle  $\mathcal{S}^c(M)$ . In general, the spinor  $\sigma$  is not necessarily parallel. In fact  $\sigma$  is parallel iff  $(M, \omega, J)$  is Kähler. However  $\sigma$  is still a harmonic spinor, a fact with several interesting applications. Here we outline a simple example. By the Lichnerowicz-Bochner formula we have

$$\nabla^* \nabla \sigma + \frac{S}{4} \sigma + \frac{1}{2} F_A \cdot \sigma = \mathcal{D}^* \mathcal{D} \sigma = 0$$

Integrating by parts gives

$$\int_M |\nabla \sigma|^2 + \frac{S}{4} + \frac{1}{2} \langle F_A \cdot \sigma, \sigma \rangle = 0.$$

By straightforward calculations one can show

$$\begin{aligned}|\nabla \sigma|^2 &= \frac{1}{16} |\nabla J|^2 \\ \langle F_A \cdot \sigma, \sigma \rangle &= -\frac{2\pi}{(m-1)!} C_1 \wedge \omega^{m-1}\end{aligned}$$

where  $C_1$  is the first Chern form. Therefore we get the following interesting formula due to Blair [Bl92]

$$\int_M \left( \frac{1}{4} |\nabla J|^2 + S \right) \frac{\omega^m}{m!} = 4\pi \int_M C_1 \wedge \frac{\omega^{m-1}}{(m-1)!}$$

For more substantial applications of the  $\text{spin}^c$  structure in symplectic geometry we refer to the work of Taubes [Ta94, Ta95].

We now choose our orthonormal basis  $e_1, \dots, e_{2m}$  so that  $e_{m+i} = J e_i$ . By a slight abuse of notation, we denote  $e_{\bar{i}} = e_{m+i} = J e_i$ . And similarly the index  $\bar{i}$  denote  $m+i$ . Hence, with  $\sigma = 1$  being the parallel  $\text{spin}^c$  spinor, we have

$$\langle \sigma, e_i e_j \cdot \sigma \rangle = -\delta_{ij}, \quad \langle \sigma, e_{\bar{i}} e_{\bar{j}} \cdot \sigma \rangle = -\delta_{ij}, \quad (2.7)$$

$$\langle \sigma, e_{\bar{i}} e_j \cdot \sigma \rangle = -\sqrt{-1} \delta_{ij}, \quad \langle \sigma, e_i e_{\bar{j}} \cdot \sigma \rangle = \sqrt{-1} \delta_{ij}. \quad (2.8)$$

Now we compute

$$\begin{aligned}
\langle \Phi(R^L(h)), \Phi(h) \rangle &= \sum_{i,j,k,l,p=1}^{2m} F_{jp} h_{ip} h_{kl} \langle e_i \cdot \sigma \otimes e^j, e_k \cdot \sigma \otimes e^l \rangle \\
&= - \sum_{i,j,k,p=1}^{2m} F_{jp} h_{ip} h_{kj} \langle \sigma, e_i e_k \cdot \sigma \rangle \\
&= \sum_{j,p=1}^{2m} \sum_{i=1}^m F_{jp} h_{ip} h_{ij} - \sum_{j,p=1}^{2m} \sum_{i=1}^m F_{jp} h_{ip} h_{\bar{i}j} (\sqrt{-1}) \\
&\quad - \sum_{j,p=1}^{2m} \sum_{i=1}^m F_{jp} h_{\bar{i}p} h_{ij} (-\sqrt{-1}) + \sum_{j,p=1}^{2m} \sum_{i=1}^m F_{jp} h_{\bar{i}p} h_{\bar{i}j}.
\end{aligned}$$

As the curvature of a unitary connection on a line bundle,  $F$  is purely imaginary. Hence taking the real part (and using the skew symmetry) yields:

$$\operatorname{Re} \langle \Phi(R^L(h)), \Phi(h) \rangle = -2\sqrt{-1} \sum_{j,p=1}^{2m} \sum_{i=1}^m F_{jp} h_{ip} h_{\bar{i}j}. \quad (2.9)$$

Similarly,

$$\operatorname{Re} \langle \Phi(\operatorname{Ric}(h)), \Phi(h) \rangle = \sum_{i,j,p=1}^{2m} R_{ip} h_{pj} h_{ij}. \quad (2.10)$$

We are now ready to prove

**Theorem 2.4** *If  $(M, g_0)$  is a compact Kähler manifold with nonpositive Ricci curvature, then  $\nabla^* \nabla h - 2\mathring{R}h$  is positive semi-definite on  $S^2(M)$ . That is,*

$$\langle \nabla^* \nabla h - 2\mathring{R}h, h \rangle \geq \langle \mathcal{D}\Phi(h), \mathcal{D}\Phi(h) \rangle \geq 0,$$

for any  $h \in S^2(M)$ . Moreover, in the case of negative Ricci curvature,  $\nabla^* \nabla h - 2\mathring{R}h = 0$  iff  $\mathcal{D}\Phi(h) = 0$  and  $h$  is skew-hermitian.

**Proof:** Since  $\langle \mathcal{D}\Phi(h), \mathcal{D}\Phi(h) \rangle \geq 0$ , we have, by Lemmas 2.3 and 2.2,

$$\begin{aligned}
0 &\leq \langle \mathcal{D}^* \mathcal{D}\Phi(h), \Phi(h) \rangle \\
&= \operatorname{Re} \langle \Phi(\nabla^* \nabla h - 2\mathring{R}h + R^L(h) + \operatorname{Ric}(h)), \Phi(h) \rangle \\
&= \langle \nabla^* \nabla h - 2\mathring{R}h, h \rangle + \operatorname{Re} \langle \Phi(R^L(h)), \Phi(h) \rangle + \operatorname{Re} \langle \Phi(\operatorname{Ric}(h)), \Phi(h) \rangle.
\end{aligned}$$

That is,

$$\langle \nabla^* \nabla h - 2\mathring{R}h, h \rangle = \langle \mathcal{D}\Phi(h), \mathcal{D}\Phi(h) \rangle - [\operatorname{Re} \langle \Phi(\operatorname{Ric}(h)), \Phi(h) \rangle + \operatorname{Re} \langle \Phi(R^L(h)), \Phi(h) \rangle].$$

For  $L = K^{-1}$ , as we remarked earlier, the curvature form  $F = -\sqrt{-1}\rho$  where  $\rho$  is the Ricci form. Since  $g$  is Kähler, there is an orthonormal basis  $e_1, \dots, e_{2m}$  such that  $e_{m+i} = Je_i$  ( $1 \leq i \leq m$ ) and the Ricci curvature is diagonal in this basis, i.e.  $R_{ij} c_i \delta_{ij}$  ( $1 \leq i, j \leq 2m$ ) with  $c_i = c_{m+i}$ . Now

$$\rho(e_i, e_j) = \operatorname{Ric}(Je_i, e_j) = \begin{cases} c_i \delta_{m+i,j}, & 1 \leq i \leq m \\ -c_i \delta_{m-i,j}, & m+1 \leq i \leq 2m \end{cases}.$$

It follows then from (2.9) that

$$\operatorname{Re}\langle\Phi(R^L(h)),\Phi(h)\rangle-2\sum_{j=1}^m\sum_{i=1}^mc_j(h_{i\bar{j}}h_{\bar{i}j}-h_{ij}h_{\bar{i}\bar{j}}).$$

By (2.10),

$$\operatorname{Re}\langle\Phi(\operatorname{Ric}(h)),\Phi(h)\rangle=\sum_{i,j=1}^{2m}c_ih_{ij}^2.$$

Hence,

$$\langle\nabla^*\nabla h-2\overset{\circ}{R}h,h\rangle\geq-\left[\sum_{i,j=1}^{2m}c_ih_{ij}^2-2\sum_{j=1}^m\sum_{i=1}^mc_i(h_{i\bar{j}}h_{\bar{i}j}-h_{ij}h_{\bar{i}\bar{j}})\right].$$

When  $c_i\leq 0$  the right hand side is nonnegative by the Cauchy-Schwarz inequality.

If  $c_i < 0$ , then  $\nabla^*\nabla h-2\overset{\circ}{R}h=0$  if and only if  $\mathcal{D}\Phi(h)=0$  and

$$h_{ij}=-h_{\bar{i}\bar{j}},\quad h_{i\bar{j}}=h_{\bar{i}j}.$$

That is,  $h(JX,JY)=-h(X,Y)$ . It follows that  $h$  is skew Hermitian. ■

Similarly, we have

**Theorem 2.5** *If  $(M,g_0)$  is a compact Kähler manifold with nonnegative Ricci curvature, then the Lichnerowicz Laplacian is positive semi-definite on  $S^2(M)$ . That is,*

$$\langle\mathcal{L}_gh,h\rangle\geq\langle\mathcal{D}\Phi(h),\mathcal{D}\Phi(h)\rangle\geq 0,$$

for any  $h\in S^2(M)$ . Moreover, in the case positive Ricci curvature,  $\mathcal{L}_gh=0$  if and only if  $\mathcal{D}\Phi(h)=0$  and  $h$  is Hermitian.

**Proof:** The Lichnerowicz Laplacian is

$$\mathcal{L}_gh=\nabla^*\nabla h-2\overset{\circ}{R}h+2\operatorname{Ric}\circ h.$$

The same proof as above now goes through for  $\mathcal{L}_g$ . ■

### 3 Stability of Kähler-Einstein metrics

Theorem 2.4 says that for a Kähler-Einstein  $(N,g_0)$  with non-positive scalar curvature the operator

$$\nabla^*\nabla h-2\overset{\circ}{R}h \tag{3.1}$$

is semi-positive definite on symmetric 2-tensors. A natural and important question is to identify the kernel space

$$W_{g_0}=\{h\mid\operatorname{tr}_{g_0}h=0,\delta h=0,\nabla^*\nabla h-2\overset{\circ}{R}h=0\} \tag{3.2}$$

on the space of transverse traceless symmetric 2-tensors. The case  $c=0$  is essentially a Calabi-Yau manifold, which has been studied with other manifolds admitting parallel spinors in our previous paper [DWW04]. We now focus on the case  $c<0$ .

By the proof of Theorem 2.4

$$W_{g_0}=\{h\mid\operatorname{tr}_{g_0}h=0,\delta h=0,h(J,J)=-h,\mathcal{D}\Phi(h)=0\} \tag{3.3}$$

As before we choose our orthonormal basis  $e_1, \dots, e_{2m}$  so that  $e_{m+i} = Je_i$ .

Now for  $1 \leq i \leq m$  set

$$X_i = \frac{e_i - \sqrt{-1}Je_i}{\sqrt{2}}, \quad \bar{X}_i = \frac{e_i + \sqrt{-1}Je_i}{\sqrt{2}}.$$

Then  $\{X_1, \dots, X_m\}$  is a local unitary frame for  $T^{1,0}M$  and let its dual frame be  $\{\theta^1, \dots, \theta^m\}$ . As  $h$  is skew-Hermitian we have

$$h(X_i, \bar{X}_j) = h(\bar{X}_j, X_i) = 0 \quad (3.4)$$

By straightforward computation we have for  $h \in W = W_{g_0}$ ,

$$\Phi(h) = h(\bar{X}_i, \bar{X}_j)\bar{\theta}^i \otimes \bar{\theta}^j. \quad (3.5)$$

This can be identified with

$$\Psi(h) = h(\bar{X}_i, \bar{X}_j)\bar{\theta}^i \otimes X_j \in \wedge^{0,1}(\Theta), \quad (3.6)$$

where  $\Theta$  is the holomorphic tangent bundle. We compute

$$\begin{aligned} \mathcal{D}\Phi(h) &= \sum_{k=1}^m (\nabla_{e_k} h(\bar{X}_i, \bar{X}_j)e_k \cdot \bar{\theta}^i \otimes \bar{\theta}^j + \nabla_{e_{\bar{k}}} h(\bar{X}_i, \bar{X}_j)e_{\bar{k}} \cdot \bar{\theta}^i \otimes \bar{\theta}^j) \\ &= \sum_{k=1}^m (\nabla_{e_k} h(\bar{X}_i, \bar{X}_j)(\bar{\theta}^k \wedge \bar{\theta}^i - \delta_{ik}) \otimes \bar{\theta}^j + \nabla_{e_{\bar{k}}} h(\bar{X}_i, \bar{X}_j)(\sqrt{-1}\bar{\theta}^k \wedge \bar{\theta}^i - \sqrt{-1}d_{ik}) \otimes \bar{\theta}^j) \\ &= \sqrt{2} \sum_{k=1}^m (\nabla_{\bar{X}_k} h(\bar{X}_i, \bar{X}_j)\bar{\theta}^k \wedge \bar{\theta}^i \otimes \bar{\theta}^j - \nabla_{X_k} h(\bar{X}_i, \bar{X}_j)\bar{\theta}^j) \end{aligned}$$

With  $\Phi(h)$  identified as  $\Psi(h) \in \wedge^{0,1}(\Theta)$ , the above calculation shows that the Dirac operator is then  $\sqrt{2}(\bar{\partial} - \bar{\partial}^*)$  on  $\Psi(h)$ . Therefore  $\mathcal{D}\Phi(h) = 0$  iff  $\Phi(h)$  is harmonic. On the other hand

$$\begin{aligned} \delta h &= \sum_{k=1}^m (\nabla_{e_k} h(e_k, \cdot) + \nabla_{e_{\bar{k}}} h(e_{\bar{k}}, \cdot)) \\ &= \sum_{k=1}^m (\nabla_{X_k} h(\bar{X}_k, \cdot) + \nabla_{\bar{X}_k} h(X_k, \cdot)) \\ &= \sum_{j,k=1}^m (\nabla_{X_k} h(\bar{X}_k, \bar{X}_j)\bar{\theta}^j + \nabla_{\bar{X}_k} h(X_k, X_j)\theta^j) \end{aligned}$$

where in the last step we used (3.4). This shows that  $\delta h = 0$  automatically holds if  $\Psi(h)$  is harmonic. Therefore we have an injective homomorphism

$$\Psi : W_{g_0} \rightarrow H^1(N, \Theta).$$

The image obviously consists of symmetric infinitesimal complex deformations. To show that  $\Psi$  is in fact onto we need to show all infinitesimal complex deformations are symmetric. For this purpose we need a digression.

Let  $N$  be a Kähler manifold with Kähler metric  $\omega = \sqrt{-1}g_{i\bar{j}}dz^i \wedge d\bar{z}^j$ . Given  $\Psi = a_j^i d\bar{z}^j \otimes \frac{\partial}{\partial z^i} \in \wedge^{0,1}(T^{1,0}N)$ , we can consider the  $(0,2)$ -form

$$\psi = g_{i\bar{l}} a_j^i d\bar{z}^j \wedge d\bar{z}^l.$$

**Remark** One can work with a local unitary frame  $\{X_1, \dots, X_m\}$  and its dual frames equally well, but it seems the calculations are easier working with local coordinates.

We calculate

$$\begin{aligned}\bar{\partial}\Psi &= \frac{\partial a_l^i}{\partial \bar{z}^j} d\bar{z}^j \wedge dz^l \otimes \frac{\partial}{\partial z^i} \\ \bar{\partial}^* \Psi &= -g^{k\bar{l}} \frac{\partial}{\partial \bar{z}^l} \lrcorner \nabla_{\frac{\partial}{\partial z^k}} \Psi \\ &= -g^{k\bar{l}} \left( \frac{\partial a_l^i}{\partial z^k} + \Gamma_{kp}^i a_l^p \right) \frac{\partial}{\partial z^i}\end{aligned}$$

Suppose now that  $\Psi$  is harmonic, i.e.  $\bar{\partial}\Psi = 0, \bar{\partial}^* \Psi = 0$ . Then we have

$$\frac{\partial a_l^i}{\partial \bar{z}^j} = \frac{\partial a_j^i}{\partial \bar{z}^l} \quad (3.7)$$

and

$$g^{k\bar{l}} \left( \frac{\partial a_l^i}{\partial z^k} + \Gamma_{kp}^i a_l^p \right) = 0 \quad (3.8)$$

Thus

$$\bar{\partial}\psi = \left( g_{i\bar{l}} \frac{\partial a_j^i}{\partial \bar{z}^q} + a_j^i \frac{\partial g_{i\bar{l}}}{\partial \bar{z}^q} \right) d\bar{z}^q \wedge d\bar{z}^j \wedge dz^l = 0,$$

here we used 3.7 and the fact that  $\frac{\partial g_{i\bar{l}}}{\partial \bar{z}^q}$  is symmetric in  $l$  and  $q$ .

We calculate

$$\begin{aligned}\bar{\partial}^* \psi &= -g^{k\bar{l}} \frac{\partial}{\partial \bar{z}^l} \lrcorner \nabla_{\frac{\partial}{\partial z^k}} \psi \\ &= -g^{k\bar{l}} \frac{\partial}{\partial \bar{z}^l} \lrcorner \left[ \left( g_{i\bar{q}} \frac{\partial a_j^i}{\partial z^k} + \frac{\partial g_{i\bar{q}}}{\partial z^k} a_j^i \right) d\bar{z}^j \wedge d\bar{z}^q \right] \\ &= -g^{k\bar{l}} \left[ \left( g_{i\bar{q}} \frac{\partial a_l^i}{\partial z^k} + \frac{\partial g_{i\bar{q}}}{\partial z^k} a_l^i \right) d\bar{z}^q - \left( g_{i\bar{l}} \frac{\partial a_j^i}{\partial z^k} + \frac{\partial g_{k\bar{l}}}{\partial z^i} a_j^i \right) d\bar{z}^j \right] \\ &= -g_{i\bar{q}} g^{k\bar{l}} \left( \frac{\partial a_l^i}{\partial z^k} + \Gamma_{kp}^i a_l^p \right) d\bar{z}^q + \left( \frac{\partial a_j^i}{\partial z^i} + \frac{\partial \log \det G}{\partial z^i} a_j^i \right) d\bar{z}^j \\ &= \left( \frac{\partial a_j^i}{\partial z^i} + \frac{\partial \log \det G}{\partial z^i} a_j^i \right) d\bar{z}^j\end{aligned}$$

where in the last step we used 3.8. Therefore

$$\begin{aligned}\bar{\partial}\bar{\partial}^* \psi &\left( \frac{\partial^2 a_j^i}{\partial z^i \partial \bar{z}^l} + \frac{\partial \log \det G}{\partial z^i} \frac{\partial a_j^i}{\partial \bar{z}^l} + \frac{\partial^2 \log \det G}{\partial z^i \partial \bar{z}^l} a_j^i \right) d\bar{z}^l \wedge d\bar{z}^j \\ &= -\frac{\partial^2 \log \det G}{\partial z^i \partial \bar{z}^l} a_j^i d\bar{z}^j \wedge d\bar{z}^l\end{aligned}$$

where in the last step we used the fact that the first two coefficients are symmetric in  $j$  and  $l$  by (3.7). To summarize we have

$$\bar{\partial}\bar{\partial}^* \psi = R_{i\bar{l}} a_j^i d\bar{z}^j \wedge d\bar{z}^l \quad (3.9)$$

**Theorem 3.1** *Let  $(N, \omega_0)$  be a compact Kähler-Einstein manifold with negative scalar curvature. Suppose  $\Psi = a_j^i d\bar{z}^j \otimes \frac{\partial}{\partial z^i} \in \wedge^{0,1}(T^{1,0}N)$  is harmonic. Then  $\psi = g_{i\bar{i}} a_j^i d\bar{z}^j \wedge dz^i = 0$ , i.e.  $g_{i\bar{i}} a_j^i$  is symmetric in  $i$  and  $j$ .*

**Proof:** By the assumption we have  $R_{i\bar{j}} = c g_{i\bar{j}}$  with  $c < 0$ . By (3.9)

$$\overline{\partial\partial^*} \psi = c\psi.$$

Therefore

$$c \int_N |\psi|^2 = \int_N \langle \overline{\partial\partial^*} \psi, \psi \rangle = \int_N |\overline{\partial^*} \psi|^2$$

Since  $c < 0$  we must have  $\psi = 0$ . ■

**Remark** The discussion in Besse [Bes87](12.96) contains some mistakes. The claim that skew-symmetric infinitesimal complex deformations are in one-to-one correspondence with holomorphic 2-forms is wrong. As the above calculation shows that in general for a harmonic  $\Psi \in H^1(N, \Theta)$  the corresponding  $(0, 2)$ -form  $\psi$  is not harmonic, hence  $\bar{\psi}$  is not holomorphic. The vanishing of the space of skew-symmetric infinitesimal complex deformations on a Kähler-Einstein manifold  $N$  with negative scalar curvature has nothing to do with the Hodge number  $h^{2,0} = h^{0,2}$ . Take a compact complex hyperbolic surface  $N$ . By Calabi-Vesentini [CV60]  $H^1(N, \Theta) = 0$ . On the other hand, since the signature  $\tau(N) > 0$  and the Euler number  $\chi(N) = 3 \tau(N)$ , one can easily see by the Hodge index theorem that  $h^{2,0}(N) \neq 0$  unless  $N$  has the same Betti numbers as  $CP^2$ , then a very special example constructed by Mumford. Therefore there are compact complex hyperbolic surfaces  $N$  with  $H^1(N, \Theta) = 0$  and  $h^{2,0}(N) \neq 0$ .

**Remark** In the Ricci flat case the above calculation shows that for a harmonic  $\Psi \in H^1(N, \Theta)$  the corresponding  $(0, 2)$ -form  $\psi$  is indeed harmonic. Conversely one can show that a harmonic  $(0, 2)$ -form  $\psi$  gives rise to a skew-symmetric infinitesimal complex deformation. Therefore the space of skew-symmetric infinitesimal complex deformations can be identified as the space of holomorphic  $(2, 0)$ -forms.

As a corollary of Theorem 3.1 we now have a clear understanding of the kernel  $W_{g_0}$  of (3.2).

**Theorem 3.2** *Let  $(N, g_0, J_0)$  be a Kähler-Einstein manifold with negative scalar curvature. Then  $\Psi : W_{g_0} \rightarrow H^1(N, \Theta)$  is an isomorphism.*

By the Kodaira-Spencer theory,  $H^1(N, \Theta)$  is the space of infinitesimal complex deformations on  $N$ . In general these infinitesimal deformations may not be integrable. But if they are integrable, then the premoduli space of complex structures on  $N$  is a manifold near  $J$ , with  $H^1(N, \Theta)$  as the tangent space (see the next section for more discussion). In this case we can deduce various local results which we now explain. The argument is by now standard, see [DWW04] and Besson-Courtois-Gallot [BCG91] where same type of results are established for Einstein manifolds with negative sectional curvature.

We consider two well-known functionals, in addition to the first eigenvalue  $\lambda(g)$  considered in [DWW04]. For a compact manifold  $(M, g)$  of dimension  $n$

$$K(g) = \int_M |S_g|^{n/2} dV_g \tag{3.10}$$

and

$$Y(g) = \frac{\int_M S_g dV_g}{\text{Vol}(g)^{1-\frac{2}{n}}}. \tag{3.11}$$

**Theorem 3.3** *Let  $(N, g_0, J_0)$  be a compact Kähler-Einstein manifold with negative scalar curvature. Suppose all infinitesimal complex deformations of  $J_0$  are integrable. Then there exists a neighborhood  $\mathcal{U}$  of  $g_0$  in the space of smooth Riemannian metrics on  $N$  such that*

$$\forall g \in \mathcal{U} \quad K(g) \geq K(g_0)$$

*and equality holds iff  $g$  is a Kähler-Einstein metric with negative scalar curvature. Moreover all Einstein metrics in  $\mathcal{U}$  are Kähler-Einstein with negative scalar curvature.*

Since all infinitesimal complex deformations of  $J_0$  are integrable, the the premoduli space of complex structures on  $N$  is an manifold near  $J_0$ , with  $H^1(N, \Theta)$  as the tangent space. By the uniqueness of Kähler-Einstein metric with negative scalar curvature and the implicit function theorem, the moduli space  $\mathcal{E}$  of Kähler-Einstein metrics is an orbifold near  $g_0$ , with  $W_{g_0} \cong H^1(N, \Theta)$  as the tangent space.

Both functionals  $K$  and  $Y$  are scaling invariant, therefore we can restrict ourselves to the space of Riemannian metric of volume 1, denoted by  $\mathcal{M}$ . By Ebin's slice theorem, there is a real submanifold  $\mathcal{S}$  containing  $g_0$ , which is a slice for the action of the diffeomorphism group on  $\mathcal{M}$ . The tangent space

$$T_{g_0}\mathcal{S} = \{h | \delta_{g_0}h = 0, \int_N \text{tr}_{g_0} h dV_{g_0} = 0.\} \quad (3.12)$$

Let  $\mathcal{C} \subset \mathcal{S}$  be the submanifold of constant scalar curvatures metrics.

We need the following simple lemma from [BCG91]

**Lemma 3.4** *Let  $g$  be a metric with scalar curvature a negative constant and  $g'$  a metric conformal to  $g$ . Then  $K(g') \geq K(g)$  and equality holds iff  $g' = g$ .*

By this Lemma and the solution of the Yamabe problem, we only need to prove  $g_0$  is a local minimum for the functional  $K$  on  $\mathcal{C}$ . So it suffices to work on  $\mathcal{C}$ . It is easy to see

$$T_{g_0}\mathcal{C} = \{h | \delta_{g_0}h = 0, \text{tr}_{g_0}h = 0.\} \quad (3.13)$$

Restricted on  $\mathcal{C}$  and in a neighborhood of  $g_0$ , the functional  $K$  becomes  $K(g) = |S_g|^m = (-S_g)^m$ . Therefore to prove that  $g_0$  is a local minimum for  $K$  on  $\mathcal{C}$  is equivalent to prove that  $Y$  has a local maximum at  $g_0$  on  $\mathcal{C}$ . It is well known that  $g_0$  is a critical point for  $Y$  and its Hessian at  $g_0$  is given by

$$D^2Y(h, h) = -\frac{1}{2} \int_N \langle \nabla^* \nabla h - 2\overset{\circ}{R}h, h \rangle. \quad (3.14)$$

$\mathcal{C}$  contains the finite dimensional submanifold  $\mathcal{E}$  of Kähler-Einstein metrics, with tangent space  $W_{g_0} \cong H^1(N, \Theta)$ . For any  $g \in \mathcal{E}$ , let  $J$  be the associated complex structure. We have

$$\rho_g = \frac{S_g}{2m} \omega_g$$

where  $\omega_g$  is the associated Kähler form and  $\rho_g$  the Ricci form. Therefore

$$Y(g) = -4m\pi \left( \frac{-C_1(N, J)^m [N]}{m!} \right)^{\frac{1}{m}},$$

where  $C_1(N, J)$  is the first Chern class of  $(N, J)$ . Thus  $Y$  is constant on  $\mathcal{E}$ . Moreover, we have proved that  $D^2Y$  is negative definite on its normal bundle. Therefore there is a possibly smaller neighborhood of  $\mathcal{E} \subset \mathcal{C}$ , still denoted by  $\mathcal{U}$ , such that

$$\forall g \in \mathcal{U} - \mathcal{E}, \quad Y(g) < Y(g_0).$$

This proves Theorem 3.3.

We could have used  $\lambda(g)$  instead of  $Y(g)$  for the proof, as in [DWW04].

**Theorem 3.5** *Let  $(N, g_0, J_0)$  be a compact Kähler-Einstein manifold with negative scalar curvature. Suppose all infinitesimal complex deformations of  $J_0$  are integrable. Then there exists a neighborhood  $\mathcal{U}$  of  $g_0$  in the space of smooth Riemannian metrics on  $N$  such that for any metric  $g \in \mathcal{U}$  with scalar curvature  $S_g \geq S_{g_0}$*

$$\text{Vol}(N, g) \geq \text{Vol}(N, g_0)$$

*and equality holds iff  $g$  is a Kähler-Einstein metric with negative scalar curvature.*

**Remark** Though only a local result, it is quite remarkable to have volume comparison under a lower bound for the scalar curvature. The scalar curvature is a very weak geometric quantity and its effect on a general Riemannian manifold  $(M, g)$  of dimension  $n$  can only be detected infinitesimally by the following expansion for the volume of a geodesic ball  $B(p, r)$

$$\text{Vol}(B(p, r)) = \omega_n r^n \left( 1 - \frac{S_g(p)}{6(n+2)} r^2 + O(r^3) \right) \quad \text{as } r \rightarrow 0$$

**Proof:** We take the same  $\mathcal{U}$  in Theorem 3.3. Then  $\forall g \in \mathcal{U}$  with  $S_g \geq S_{g_0}$ , we have  $|S_g|^m \leq |S_{g_0}|^m$  since  $S_g < 0$ . Therefore  $K(g) \leq |S_{g_0}|^m \text{Vol}(N, g)$  while  $K(g_0) = |S_{g_0}|^m \text{Vol}(N, g_0)$ . The result then follows from Theorem 3.3. ■

Theorem 3.3 has another interesting interpretation. Recall that the Yamabe invariant of a compact Riemannian manifold  $(M, g)$  of dimension  $n$  is

$$\mu(g) = \inf_{f \in C^\infty(M), f > 0} Y(f^{\frac{n-2}{4}} g) \quad (3.15)$$

and it is a conformal invariant. The Yamabe number of  $M$  is defined as

$$\sigma(M) = \sup_g \mu(g). \quad (3.16)$$

We can now reformulate Theorem 3.3 as follows

**Theorem 3.6** *Let  $(N, g_0, J_0)$  be a compact Kähler-Einstein manifold with negative scalar curvature. Suppose all infinitesimal complex deformations of  $J_0$  are integrable. Then  $g_0$  is a local maximum of the Yamabe invariant.*

We end with a few remarks. In [Sch89] Schoen made the following conjecture

**Conjecture 3.7** *Let  $(M, g_0)$  be a compact hyperbolic manifold. Then  $\sigma(M)$  is achieved by  $g_0$  and only by  $g_0$ . In other words*

$$\forall g \quad \mu(g) \leq \mu(g_0)$$

*and equality holds iff  $g$  is conformal to  $g_0$ .*

It is also reasonable to make the same conjecture for other compact locally symmetric spaces with negative sectional curvature. In complex dimension 2 there have been some remarkable results proved by LeBrun using Seiberg-Witten invariants. For example he proved that on any compact Kähler-Einstein surface with negative scalar curvature the Yamabe number is achieved by the Kähler-Einstein metric. See [Le95, Le99].

In view of LeBrun's result and Theorem 3.6 it is tempting to extend the conjecture to Kähler-Einstein manifolds  $(N, g_0)$  with negative scalar curvature in higher dimensions, namely that  $g_0$  should be a global maximum of the Yamabe invariant provided all its infinitesimal complex deformations are integrable. But in general this is not true. In fact any compact and simply connected manifold  $M$  of dimension  $\geq 5$  has  $\sigma(M) \geq 0$ . This is trivial if  $\Sigma$  admits a metric of positive scalar curvature. Otherwise it is proved by Petean [Pe00].

Another intriguing question is whether Theorem 3.3 and Theorem 3.6 are still true if there are non-integrable complex deformations. Then we have infinitesimal Kähler-Einstein deformations which can not be integrated to Kähler-Einstein metrics, but they may be integrated to Einstein metrics which are not Kähler.

## 4 Kodaira-Spencer theory

The deformation theory of complex structures was introduced by Kodaira-Spencer in their seminal work [KS58-1, KS58-2, KS60]. This deep theory has played and still plays significant role in the theory of complex manifolds. The relation between Kodaira-Spencer theory and the deformation of Einstein metrics has been studied by Koiso [Ko83]. We review some relevant facts here in this section and discuss some examples in more detail.

Let  $M$  be a compact complex manifold and  $\Theta$  the (sheaf of germs of the) holomorphic tangent bundle of  $M$ . According to the Kodaira-Spencer theory, the infinitesimal complex deformations are described by the cohomology group  $H^1(M, \Theta)$ . For our purpose, we are interested in the integrability of infinitesimal complex deformations. Namely, when does every infinitesimal deformation actually arise from a deformation of complex structures? Let's recall first the so-called Theorem of Existence in the Kodaira-Spencer theory [Kod86, Theorem 5.6].

**Theorem 4.1 (Kodaira-Spencer)** *Let  $M$  be a compact complex manifold. If  $H^2(M, \Theta) = 0$ , then there is a complex analytic family with base  $B$ ,  $0 \in B \subset \mathbb{C}^m$ , such that the fiber at 0 is  $M$  and the Kodaira-Spencer map at 0 is an isomorphism from  $T_0B$  onto  $H^1(M, \Theta)$ .*

Recall that the Kodaira-Spencer map for a differentiable family of compact complex manifolds assigns a tangent vector of the base to the infinitesimal deformation along that direction. Thus the condition  $H^2(M, \Theta) = 0$  implies that all infinitesimal complex deformations are integrable.

We now discuss some examples from [Kod86], where the reader is referred to for complete detail.

**Example: 1).** Blowups of  $\mathbb{C}\mathbb{P}^2$ . For  $M = \mathbb{C}\mathbb{P}^2 \#_k \overline{\mathbb{C}\mathbb{P}^2}$  and  $k \geq 5$ , one has  $H^0(M, \Theta) = H^2(M, \Theta) = 0$ . Hence, all infinitesimal deformations are integrable. Incidentally, for  $k \leq 4$ ,  $H^1(M, \Theta) = 0$ . Therefore the complex structure is rigid in these cases.

It is well-known that there exists Kähler-Einstein metrics on  $M$  if and only if  $3 \leq k \leq 8$  by Tian's work [T97]. However, these Kähler-Einstein metrics have positive scalar curvature. Hence our results do not apply.

**Example: 2).** Surfaces of arbitrary degree. For a non-singular surface  $M$  of degree  $h$  in  $\mathbb{C}\mathbb{P}^3$ , one has

$$\dim H^2(M, \Theta) = \frac{1}{2}(h-2)(h-3)(h-5).$$

Thus,  $H^2(M, \Theta) = 0$  for  $h = 2, 3, 5$ . Since  $c_1(M) = (4-h)H$  where  $H$  is the hyperplane class,  $M$  has Kähler-Einstein metrics with negative scalar curvature if  $h \geq 5$ , by the Calabi-Aubin-Yau Theorem. Hence our results apply to the non-singular surface of degree 5.

As one can see here, in general, the condition  $H^2(M, \Theta) = 0$ , which guarantees the integrability of all infinitesimal complex deformations, is very restrictive. Indeed, there are many examples which do not satisfy this condition but still, all their infinitesimal complex deformations are integrable. In fact, understanding the reason behind this is one of the motivations for Kodaira-Spencer.

From the Kodaira-Spencer theory, if an infinitesimal complex deformation  $\theta \in H^1(M, \Theta)$  is integrable, then

$$[\theta, \theta] = 0.$$

This is in fact the first order obstruction. One thus expects that there should be non-integrable infinitesimal deformations. However, the cohomology group  $H^1(M, \Theta)$  turns out to be surprisingly difficult to compute. And in the many examples where it can be computed, the infinitesimal deformations turn out to be integrable.

Recall that a complex analytic family of compact complex manifolds is said to be effective (or minimal) if its Kodaira-Spencer map is injective. It is called a complete (or versal) family if every other (sufficiently small) family of deformations can be induced from this family via pullback of a

holomorphic map. Now whenever there is an effective complete family with base  $B \ni 0$  a domain in  $\mathbb{C}^m$ , such that the fiber at 0 is  $M$ , Kodaira-Spencer defines the number of moduli  $m(M) = m$  to be the dimension of the base. Then the question of whether all infinitesimal deformations are integrable can be reinterpreted as when the equality  $m(M) = \dim H^1(M, \Theta)$  holds, of which Kodaira-Spencer refers as the fundamental guiding question in the Kodaira-Spencer theory.

By the Theorem of Completeness in the Kodaira-Spencer theory [Kod86, Theorem 6.1], which says that a complex analytic family of compact complex manifolds with surjective Kodaira-Spencer map is complete, the complex analytic family in Theorem 4.1 is an effective complete family with base dimension  $\dim H^1(M, \Theta)$ .

More generally, if  $M$  is a compact complex manifold for which there is a complex analytic family of deformations whose Kodaira-Spencer map is surjective, then all infinitesimal deformations are integrable. This is the case where many examples can be found. (We will see the converse in a moment.)

**Example: 3).** Hypersurfaces in  $\mathbb{C}\mathbb{P}^n$ . If  $M$  is a hypersurface in  $\mathbb{C}\mathbb{P}^n$  of degree  $d$ , one can construct a complex analytic family of deformations of  $M$  by varying the coefficients of the defining equation of  $M$ . This family has surjective Kodaira-Spencer map. There are many examples in this class which admits Kähler-Einstein metrics with negative scalar curvature. Let  $N \subset \mathbb{C}\mathbb{P}^{m+1}$  be a smooth algebraic hypersurface of degree  $d > m + 2$ . Then the 1st Chern class  $c_1(N) < 0$ . By the theorem of Calabi-Aubin-Yau, there is a Kähler-Einstein metric with negative scalar curvature  $g_0$  on  $N$ . It is shown in [Kod86, p219] that

$$\dim H^1(N, \Theta) = \binom{m+1+d}{d} - (m+2)^2.$$

Going back to the question of existence, without any assumptions, there is the Kuranishi Theorem [Ku62].

**Theorem 4.2 (Kuranishi)** *For any compact complex manifold  $M$ , there exists a complete complex analytic family with base  $B$ ,  $0 \in B$  such that the fiber at 0 is  $M$ . Moreover  $B$  is a complex analytic subset of  $\mathbb{C}^m$ , where  $m = \dim H^1(M, \Theta)$ , defined by  $l$  holomorphic equations, with  $l = \dim H^2(M, \Theta)$ .*

It can be deduced from Kuranishi's theorem that if every infinitesimal complex deformations are integrable, then the Kuranishi family above is a complex analytic family whose Kodaira-Spencer map is an isomorphism (and hence surjective). Thus, our integrability assumption is equivalent to the existence of complex analytic family of deformations whose Kodaira-Spencer map is surjective.

As pointed out by Koiso [Ko83], the existence of such a family has significant implication for the moduli space of (Kähler-)Einstein metrics.

**Theorem 4.3** *Let  $M$  be a compact complex manifold and  $(g, J)$  be a Kähler-Einstein structure on  $M$ . Assume that the complex structure  $J$  belongs to a complex analytic family of complex structures with surjective Kodaira-Spencer map. Moreover, if the scalar curvature is positive, assume further that there is no nonzero hermitian infinitesimal Einstein deformations and also no nonzero holomorphic vector field. Then the local premoduli space of Einstein metrics around  $g$  is a manifold with tangent space at  $g$  the space of infinitesimal Einstein deformations. Moreover, any Einstein metric in it is Kähler (with respect to some complex structure).*

It should be pointed out that there are indeed examples of compact complex manifolds with non-integrable infinitesimal complex deformations [Kod86, p319]. However, we do not know any examples of Kähler-Einstein manifolds with negative scalar curvature which does not satisfy the integrability condition. In view of the Bogmolov-Tian-Todorov theorem [Bo78, T86, To89] in the Calabi-Yau case, we have the following very interesting question.

**Question:** Is it true that on any compact Kähler-Einstein manifolds with negative scalar curvature, the universal deformation space of complex structures is smooth?

We hope to come back to this question later.

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